Chebyshev Based Spectral Representations of Neutron-Star Equations of State

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Introduction

- The equation of state describes the relationship between pressure (P) and energy density (ϵ) ;
- Given an equation of state $\epsilon = \epsilon(P)$, Oppenheimer and Volkoff solved the Einstein equations to get

$$\frac{dm}{dr} = 4\pi r^2 \epsilon \tag{1}$$

$$\frac{dP}{dr} = -(\epsilon + P)\frac{m + 4\pi r^3 P}{r(r - 2m)};$$
(2)

- Neutron star densities surpass lab limits, with no consensus on models, but mass (M) and radius (R) data can help determine their equations of state [1];
- Aim to represent the equations of state numerically with spectral methods to find the best-fit equations of state with M and R data [2];
- Simple power-law bases have been used in previous studies but show slow convergence when adding more spectral parameters for first-order and second-order phase transitions [3];
- Chebyshev polynomials can serve as a faster-converging orthogonal basis, defined by:

$$T_n(x) = \cos(n \arccos(x)) \tag{3}$$

for $n = 0, 1, 2, 3, \dots$



Figure 1: These curves correspond to 34 nuclear-theory models of the neutron stars' equations of state.



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Figure 2: These curves correspond to the equations of state in Figure 1, showing the relationship between mass and radius for different equations of state.

Casual Spectral Representation

• The sound speed of the baryonic fluid is defined by $v^2 = dp/d\epsilon$;

• To ensure the causality of the fluid, we define a velocity function

$$\Upsilon = \frac{c^2 - v^2}{v^2};\tag{4}$$

• We could either represent using a pressure-based or enthalpy-based approach;

• For pressure-based approach, reduce the sound speed equation to

$$\epsilon(P) = \epsilon_0 + (P - P_0) + \beta_{P_0}^P \Upsilon(P') dP'; \quad (5)$$

• Set Υ to be expressed by spectral method,

$$\Upsilon(P, \upsilon_a) = \Upsilon_0 \exp\left\{ \sum_{\substack{\Sigma \\ a=0}}^{N_{\text{parms}}-1} \upsilon_a (1+z) T_a(z) \right\}, \quad (6)$$

where v_a is a set of spectral parameters; • Define

P = P(h) and $\epsilon = \epsilon(P(h)) = \epsilon(h)$ where h is the enthalpy.

Figure 3: Modeling errors χ using pressure-based Chebyshev polynomial spectral expansions for first-order phase transitions.

tions.

Testing The Representations

• Define the error function $\chi^2 = \frac{1}{N} \sum_{i=0}^{N} [\log(\frac{\epsilon(v_a)}{\epsilon_i})]^2;$ • Using datasets from the first-order and second-order phase transitions to optimize the error function with the best-fitted sets of spectral coefficients;





Figure 4: Modeling errors χ using pressure-based Chebyshev polynomial spectral expansions for second-order phase transi-



Figure 5: The average modeling error for different nuclear-theory models between Chebyshev enthalpy-based and pressure-based representations.

Conclusion

• When N_{parms} increases, the Chebyshev polynomial representation converges; • Pressure-based representation does a better job than enthalpy-based representation; • In the future, we will explore how well the Chebyshev polynomial pressure-based representation models the equations of state by using M and R data with noise.

References

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